Low-loss single-mode hollow-core fiber with anisotropic anti-resonant elements

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Abstract: A hollow-core fiber using anisotropic anti-resonant tubes in the cladding is proposed for low loss and effectively single-mode guidance. We show that the loss performance and higher-order mode suppression is significantly improved by using symmetrically distributed anisotropic anti-resonant tubes in the cladding, elongated in the radial direction, when compared to using isotropic, i.e. circular, anti-resonant tubes. The effective single-mode guidance of the proposed fiber is achieved by enhancing the coupling between the cladding modes and higher-order-core modes by suitably engineering the anisotropic anti-resonant elements. With a silica-based fiber design aimed at 1.06 µm, we show that the loss extinction ratio between the higher-order core modes and the fundamental core mode can be more than 1000 in the range 1.0-1.65 µm, while the leakage loss of the fundamental core mode is below 15 dB/km in the same range.

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OCIS codes: (060.4005) Microstructured fibers; (060.2310) Fiber optics; (060.2400) Fiber properties; (060.2280) Fiber design and fabrication.

References and links


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1. Introduction

Light guidance in hollow-core fibers (HCFs) [1,2] has enabled new applications due to their extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates extraordinary properties compared to solid-core fibers: when light propagates in a gas-filled core instead of glass, it propagates faster and often with low dispersion, and the gas tolerates
well as single-mode operation, achieved by increasing the loss of higher-order modes (HOMs).

HC-AR fibers with both circular AR tubes [12,15,22] as well as more intricately shaped AR tubes [11,13] have been investigated. A significant loss reduction is possible with nested tubes inside the AR tubes [14,16,17,23], but this substantially increases complexity. Here we propose a simpler solution using anisotropic AR tubes, elongated along the fiber radial direction, which allows simultaneously achieving (a) an increased negative curvature in the core, (b) a node-free design, and (c) a larger distance from the core to the outer capillary. All these properties could not be achieved simultaneously in the previous cases [11–17,22]. Importantly, these properties offer a degree of freedom in the design to reduce the losses significantly, achieve low-loss broadband transmission, and effectively suppress HOMs. Numerical results for a silica-based design targeting 1.06 µm show that the leakage loss can be reduced 1–2 orders of magnitude (4 dB/km at 1.06 µm) compared to the standard HC-AR fiber with circular AR tubes. Moreover, the fiber is made effectively single-moded by suppressing HOMs resulting in an extinction loss ratio between the core HOMs and fundamental mode (FM) that is over 1000 in the 1.0–1.65 µm spectral range, while the FM in the same range experiences loss <15 dB/km. Such specs cannot be reproduced with the standard isotropic design without using complex designs with multiple nested AR tubes [17].

![Fig. 1. Geometry of the considered HC-AR fibers, keeping fixed the core radius \( R = 15 \) µm and silica strut thickness \( t = 0.42 \) µm. The structural parameters shown in the figure are those that optimized the leakage loss at 1.06 µm. The figures are scaled to indicate their relative size.](image)

### 2. Numerical results

Figure 1(a)-(c) shows the considered HC-AR fiber geometries, using a thick outer capillary with AR tubes on the inner wall. Design (a) is the usual case with touching isotropic (circular) AR cladding tubes. Starting from (a), the design is optimized for minimal losses, resulting in the circular design (b) and elliptical design (c). The latter is the proposed anisotropic AR tube design, here an ellipse squeezed in the azimuthal direction. Other anisotropic shapes are possible. We focus on a silica fiber designed for \( \lambda = 1.06 \) µm (i.e. for high-power Yb lasers), which has 6 AR tubes (a larger number is also feasible, and below we will specifically compare 6 vs. 8 AR tubes), a fixed core radius \( R = 15 \) µm (large enough to enable high-power transmission) and silica strut thickness \( t = 0.42 \) µm (making the first high-loss resonance occur at around \( \lambda = 0.88 \) µm [16]). This choice of strut thickness implies that light at 1.06 µm is guided in the fundamental AR transmission band, which compared to the next higher-order AR transmission bands is favorable since it performs better in terms of loss and transmission bandwidth. The ellipticity is defined as \( \eta = r_y/r_x \), where \( r_y \) is the radius in the azimuthal direction and \( r_x \) is the radius in the radial direction; in the following we keep \( r_x = 15 \) µm fixed, and \( \eta < 1 \) will reduce the loss. We used a quarter of the geometry for the numerical calculations because of mode symmetry [24], except for the bend loss calculations where a half geometry was used due to a reduced symmetry in the elliptical case. We used the same numerical method as explained in [17], which briefly explained relies on finite-element simulations to calculate the fiber modes and their propagation constants.
2.1 Optimization the leakage loss of the HC-ARFs

First the HC-AR fibers were optimized to get the lowest loss at 1.06 µm by adjusting the size of the AR elements with the core size fixed (see [17] for details on the calculations). Figure 2(a) shows the leakage loss (or confinement loss, \( \alpha_c \)) as a function of AR air-hole radius for the circular case. When the AR tube radius decreases from \( r = 15 \) to 10.2 µm, the leakage loss decreases to a minimum value of 30 dB/km, i.e. improved with around one order of magnitude. Interestingly, the AR tubes are here much smaller than the “non-touching” node-free design (i.e. where the AR tubes are reduced just enough to prevent them from touching each other), which otherwise has been considered optimal [16,22]; the reasoning has been that the core FM no longer has coupling loss to the cladding modes that in the touching-case reside in the glass intersections in these nodes. However, this would imply a sharp drop in loss as \( r \) is taken below 15 µm to the non-touching value. Instead it is continuous, indicating that the coupling loss to the cladding modes in the glass nodes is not dominating for the chosen fiber design (however, this does not mean that for other designs it is not important), and the loss instead drops smoothly because the FM phase-mismatch to the cladding modes increases gradually. This is eventually balanced with the increased loss of the FM as its evanescent tail overlaps more with the outer capillary wall as the circles shrink. This is because we here fix the core size, so the shrinking circles imply that the outer capillary becomes closer to the core.

![Figure 2(a)](image)

Fig. 2. Calculated leakage loss at 1.06 µm as a function of (a) air-hole radius for circular AR tubes and (b) ellipticity of the AR tubes keeping \( r_x = 15 \) µm. Inset: FM field profiles at 1.06 µm. (c) Loss vs. wavelength for different HC-AR fibers (dashed line: \( \lambda = 1.06 \) µm). All structures have the same core radius \( R = 15 \) µm and uniform silica strut thickness \( t = 0.42 \) µm.

Figure 2(b) shows the leakage loss for the elliptical case. As we fix the core size and decrease \( \eta \), the major axis (in radial direction) is fixed at \( r_x = 15 \) µm, and the minor axis (in the azimuthal direction) changes from \( r_y = 15 \) to 9 µm. The lowest leakage loss of 4 dB/km was obtained for \( \eta = 0.65 \), i.e., \( r_x = 15 \) µm and \( r_y = 9.80 \) µm. Thus, orders of magnitude improvement is realized by squeezing the azimuthal axis of the AR tubes. The minimum has a different explanation than the circular case, because we are here able to fix the distance from the core to the outer capillary. Instead the loss improvement obtained for \( \eta < 1 \) is due to an increased phase mismatch between the FM and cladding modes is eventually balanced by an increased leakage loss as the FM starts leaking into the voids between the ever slimmer AR ellipses.
Figure 2(c) depicts the spectral loss distribution for different HC-AR fibers. Curves 1-3 show the circular cases: case 1 (green) where the tube walls are touching each other, thus forming glass nodes in the cladding, and case 2 when the air-hole radius is reduced to 14 µm so the AR tubes no longer touch each other (black). Case 3 (blue) shows the additional reduction of leakage loss until the minimum is reached, achieved as mentioned above by shrinking the circular tubes further (thus separating them further). Finally, case 4 (red) shows the elliptical design, displaying significantly improved loss performance; the leakage loss is over one order of magnitude lower than for the best circular design. Moreover, the low-loss range spans nearly the entire near-IR, which is promising for broadband ultrafast applications.

2.2 Scaling the strut thickness

One of the limiting factors in high-power beam delivery is the so-called fraction of power in silica (FOPS), i.e., the optical power overlap with the silica cladding, which should be kept small. Figure 3(b) shows that FOPS can be reduced by reducing the strut thickness of silica (~2x10⁻⁵ at 1.06 µm for t = 0.35 µm), which is several orders of magnitude lower than the HC-PBG [4]. This makes HC-AR fibers an ideal medium to explore propagation of high power beam delivery. When comparing FOPS for circles and ellipses the performance is similar, but the advantage in using ellipses comes in terms of leakage loss, see Fig. 3(a), which has a local minimum at a wavelength controlled by the strut thickness. As the strut thickness is reduced, this minimum shifts towards lower wavelengths, but at the wavelength of 1.06 µm that the design is intended for the leakage loss becomes quite high as the strut thickness is reduced. In the elliptical case the wavelength loss-variation is instead much more flat, so the leakage loss at 1.06 µm varies only little when the strut thickness is reduced. Figure 3(c)-3(d) summarizes these trends by showing the leakage loss and FOPS vs. strut thickness at 1.06 µm.

Fig. 3. Calculated (a) leakage loss and (b) fraction of power in silica (FOPS) vs. wavelength for different strut thicknesses. Solid lines and dashed lines are calculated for t = 0.42 µm and t = 0.35 µm respectively; (c) Leakage loss and (d) FOPS vs strut thickness at 1.06 µm.
2.3 Effectively single-mode operation

HC-AR fibers with large cores are not single-moded, but they can be made effectively single-moded by engineering the shape and size of the AR tubes so the HOMs experience more loss than the FM [17]. Figure 4(a) shows the relative effective indices ($\Delta n_{\text{eff}} = n_{\text{eff}}-1$) of the first three core modes (here denoted LP$_{01}$, LP$_{11}$, and LP$_{21}$) and the first three cladding modes. The core FM (LP$_{01}$) has the highest $\Delta n_{\text{eff}}$, which remains constant as a function of ellipticity. The first three cladding modes have only slightly larger $\Delta n_{\text{eff}}$ than the first core HOM (LP$_{11}$) because the core area is only a few times larger than the area of a single cladding tube. Thus, the first HOM is located within the domain of the cladding modes, which increases phase matching to cladding modes [25,26]. This effect is more evident for the strongly elliptical AR tubes ($\eta \sim 0.60-0.70$), where the cladding modes are better phase-matched to the HOMs than to the FM, which effectively suppresses the guidance of the HOMs due to higher losses.

The FM loss decreases much more than the HOM loss when the ellipticity is decreased, and in the $\eta = 0.60-0.70$ range the HOM losses even start increasing while the FM loss remains at 4 dB/km. This shows how the HOM losses can be made higher by suitably choosing the ellipticity. The aim is maximizing the so-called HOM extinction ratio (HOMER), defined as the ratio between the loss of the HOM with the lowest loss and the FM loss [16]. The maximum HOMER was found to be ~2500 at $\eta = 0.61$, while at $\eta = 0.65$, where the lowest FM loss was found at 1.06 µm, a HOMER of ~200 is found; both high enough to make the fibers effectively single-moded. Remarkably, the $\eta = 0.61$ case is only slightly more lossy (5 dB/km) than the $\eta = 0.65$ case, so the loss penalty of maximizing HOMER is small. This again shows the design freedom of the anisotropic AR elements. The spectral loss and HOMER is shown in Fig. 4(b). The HOMER for the design with the lowest loss at 1.06 µm ($\eta = 0.65$) can be made in excess of 150 between $\lambda = 0.95-1.8$ µm, while keeping $\alpha_c < 15$ dB/km. Thus, this fiber has low-loss and is effectively single-moded over an octave of bandwidth.

Interestingly, when the loss is increased slightly to maximize HOMER ($\eta = 0.61$), HOMER>1000 between $\lambda = 1.0-1.75$ µm with $\alpha_c < 15$ dB/km from $\lambda = 1.0-1.65$ µm. Figure 5(a) shows HOMER vs. wavelength, and confirms that the elliptical case outperforms the circular case in the entire wavelength regime 0.9-2 µm. Figure 5(b) shows the bend loss ($\alpha_b$) of the considered structures, calculated in both the x and y directions. The elliptical case shows an azimuthal variation of the bend loss, evidenced by a loss peak seen only in the x-direction for low bending radii due to increased core-cladding mode coupling. Therefore the
circular case shows better bend loss performance for low bend radii for the 6-tube structure studied here.

Fig. 5. (a) Wavelength dependence of HOMER for circular (r = 10.2µm) and elliptical (η = 0.61) AR tubes (b) Bend loss vs bend radius for circular (r = 10.2µm) and elliptical (η = 0.65) AR tubes with t = 0.42 µm. The FM profiles are shown in the right hand side for a 10 cm bending radius.

Fig. 6. (a) Loss vs. wavelength (b) HOMER (c-d) Bending loss vs. bending radius for different HC-AR fibers. All structures have the same core radius R = 15 µm and uniform silica strut thickness t = 0.42 µm. All fiber designs are optimized at 1.06 µm to give minimum leakage loss. The contour plots of the fundamental air-core mode distribution are shown in the right hand side for a 10 cm bending radius. The color of the frame corresponds to the color of the line in the plot.

2.4 Comparison with other design cases

In Fig. 6 we compare the loss performance and HOMER of 6 and 8 AR tubes. The calculated loss spectra in Fig. 6(a) shows that 6 and 8 AR tubes have similar loss performance in the 0.9-1.45 µm spectral regime; the 8 tube cases have slightly lower losses, but using 6 tubes gives much broader low-loss transmission window for both circular and elliptical cases. We also note that the 6 tube cases (both circular and elliptical) are very smooth while the 8 tube cases show spectral fluctuations vs. wavelength at the end of the transmission window. Similar fluctuations have previously been attributed to presence of nodes in the cladding [16], but clearly the origin here is different as there are no nodes. The reason is instead found in the fact that for 8 tubes the core mode for longer wavelengths expands its mode field diameter so that it starts interacting weakly with cladding modes found at the outer capillary wall [27]; this leads to the observed fluctuations. Figure 6(b) shows the HOMER from which we see that 6 elliptical AR tubes have much higher HOMER compared to 8 AR tubes in the whole spectral regime. Finally, Figs. 6 (c)-6(d) show the bend loss performance in the x- and y- directions, respectively. For 8 AR tubes the elliptical case has better bend loss performance than the
circular case. We believe this is because for the 8 tubes case, the AR tubes are smaller compared to the 6 tubes case. Therefore for small bend radius, 6 AR tubes have coupling between the core modes and tube modes due to the larger AR tubes whereas the coupling is reduced in 8 AR tubes because of the smaller AR tubes. Figure 6 also shows contour plots of the fundamental modes for 10 cm bending radius in which for 6 AR tubes there is a coupling between the core modes and tube modes whereas for 8 AR tubes there is no coupling between the core modes and tube modes. Therefore, choosing 6 or 8 tubes would have to be a compromise in terms of whether loss bandwidth, HOMER or bend loss is the most important feature.

We also considered introducing the ellipticity from the optimized circular design (10.2 μm AR tube radius), but instead of squeezing the ellipse azimuthally, it was elongated radially (i.e. extending the major axis) while keeping the core size fixed. This implies that the outer capillary expands its size as the ellipticity drops, which results in an improved loss performance compared to shrinking the minor axis (the case we have discussed so far). However, our calculations showed that the HOMER could not reach the same high values as found in Fig. 5(a). This implies that it is harder to reduce the phase-mismatch between the core HOMs and the cladding modes when extending the major axis.

3. Conclusion

In summary, a novel hollow-core anti-resonant fiber design has been proposed, in which the anti-resonant elements in the cladding are anisotropic in shape, in contrast to the conventional isotropic circular shape. Numerical simulations using an elliptical shape as the anisotropic cladding element showed loss performances and effective single-mode guidance that could not be achieved with isotropic (circular) cladding elements. In both cases the guiding of light in the core is based on the anti-resonances of the struts in the cladding and the inhibited coupling mechanism. However, the anisotropic shape has improved performance because it simultaneously offers: (a) strong negative curvature in the core, (b) node-free (non-touching) anti-resonant elements, and (c) larger distance from the core to the outer capillary for a given core curvature. This gives a design degree-of-freedom essential for enhancing the performance, e.g. by fixing the size of the core and outer capillary while tuning the ellipticity to minimize loss. We studied a specific case, where a silica-based fiber was optimized for the Yb-laser wavelength of 1.06 μm. The HOM extinction ratio was over 1000 in the range $\lambda = 1.0 - 1.75 \mu m$ with FM loss <5 dB/km at 1.06 μm and <15 dB/km for $\lambda = 1.0 - 1.65 \mu m$, which relied on increasing HOM loss by reducing the phase-mismatch between them and the cladding modes, while still maintaining a large phase-mismatch between the FM and the cladding modes. These properties are extremely promising for ultra-fast extreme nonlinear optics applications exploiting fiber-based light-matter interaction [3]. The proposed design is generic, irrespective of glass composition and target wavelength, and we expect it to improve almost any type of hollow-core fiber design exploiting the anti-resonant effect.